

Marine control over negative power law scaling of mass wasting events in chalk sea cliffs with implications for future recession under the UKCP09 medium emission scenario

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1 **Marine control over negative power law scaling of mass wasting events in chalk**
2 **sea cliffs with implications for future recession under the UKCP09 medium**
3 **emission scenario**

4 Jamie Gilham, John Barlow, and Roger Moore

5 Department of Geography, University of Sussex, Brighton, UK

6 **Abstract**

7 Coastal cliff erosion represents a significant geohazard for people and infrastructure.
8 Forecasting future erosion rates is therefore of critical importance to ensuring the
9 resiliency of coastal communities. We use high precision monitoring of chalk cliffs at
10 Telscombe, UK to generate monthly mass movement inventories between August
11 2016 and July 2017. Frequency-magnitude analysis of our inventories demonstrate
12 negative power law scaling over 7 orders of magnitude and, for the first time, we report
13 statistically significant correlations between significant wave height (H_s) and power law
14 scaling coefficients (r^2 values of 0.497 and 0.590 for β and s respectively). Applying
15 these relationships allows for a quantitative method to predict erosion at the site based
16 on H_s probabilities and sea level forecasts derived from the UKCP09 medium emission
17 climate model (A1B). Monte-Carlo simulations indicate a range of possible erosion
18 scenarios over 70 years (2020-2090) and we assess the impact these may have on
19 the A259 coastal road which runs proximal to the cliffs. Results indicate a small
20 acceleration in erosion compared to those based on current conditions with the most
21 likely scenario at the site being 21.7 m of cliff recession by 2090. However, low-
22 probability events can result in recession an order of magnitude higher in some
23 scenarios. In the absence of negative feedbacks, we estimate an ~11% chance that
24 the A259 will be breached by coastal erosion by 2090.

25 Keywords: Cliff recession, power law scaling, numerical modelling

26 Introduction

27 The evolution of sea cliffs has received greater attention in the research
28 literature since the turn of the century (e.g. Lee et al., 2001; Trenhaile, 2002; Dong &
29 Guzzetti, 2005; Teixeira, 2006; Marques, 2008; Lim et al., 2010; Hurst et al., 2016).
30 However, ~~decadal scale-quantitative~~ prediction of coastal erosion in response to
31 changing environmental controls ~~over decadal time scales~~ remains a major research
32 objective (Wong et al., 2014). Sea cliff environments are more dynamic than terrestrial
33 slopes due to exogenic forcing from both sub-aerial and marine processes. These, in
34 conjunction with material properties, control the evolution of sea cliffs (Emery & Kuhn,
35 1982; Lim et al., 2010). Although the mechanical properties of rock slopes are well
36 understood (e.g. Selby, 1993; Wyllie & Mah, 2004; Stead & Coggan, 2012),
37 deterministic approaches become impractical over large spatial and temporal scales
38 due to the difficulties involved in parameterising the sub-surface and in predicting the
39 variability of environmental controls (e.g. Stead & Coggan, 2012). Therefore, an
40 alternate approach is required involving empirically-based stochastic models which
41 extrapolate limited observations over spatial and temporal scales more applicable to
42 landscape evolution (Barlow et al., 2012). Most estimates of future sea cliff erosion
43 are based on change detection between historical maps and aerial photographs (e.g.
44 Moore et al., 2003a; 2003b; Dornbusch et al., 2008). However, this approach is
45 problematic in that measured recession rates are often similar to the precision of the
46 data and extrapolating rates forward in time assumes erosion will not be influenced by
47 projected changes in exogenic boundary conditions (Rosser et al., 2005; Lim et al.,
48 2010). The ability to conduct high-precision monitoring through terrestrial or airborne
49 laser scanning (TLS or ALS) and more recently through UAV photogrammetry has

50 greatly improved the precision of coastal erosion data (Slatton et al., 2007; Remondino
51 et al., 2011; Haala & Rothermel, 2012; Hugenholtz et al., 2013; Gonçalves &
52 Henriques, 2015). However, linking cliff erosion to environmental conditions remains
53 difficult as mass wasting events often occur through progressive failure such that
54 triggering events and failure do not necessarily correlate though time (Lim et al., 2010;
55 de Vilder et al., 2017).

56 Magnitude-frequency analysis is a statistical method for characterising
57 geomorphic events in space and time (Wolman & Miller, 1960; Stark & Guzzetti, 2009).
58 A substantial amount of research has been undertaken on the methods and
59 characteristics of landslide magnitude-frequency distributions (e.g. Hovius et al., 1997,
60 1998, 2000; Pelletier et al., 1997; Stark & Hovius, 2001; Guzzetti et al., 2002; Martin
61 et al., 2002; Hergarten, 2003; Malamud, 2004; Van Den Eeckhaut et al., 2007; Brunetti
62 et al., 2009; Rossi et al., 2010). Mathematically, these studies have demonstrated that
63 a negative power law best describes landslide magnitude-frequency distributions,
64 generally expressed as (Brunetti et al., 2009):

$$55 \quad f(V_R) = sV_R^{-\beta} \quad (1)$$

66 where $f(V_R)$ is the frequency density, V_R is the magnitude of a given event and s and
67 β are empirically derived scaling parameters. Many authors have noted the variability
68 of negative power law scaling parameters due to regional characteristics such as
69 structural geology, morphology, hydrology and climate (Stark & Hovius, 2001;
70 Dussauge et al., 2003; Brardinoni & Church, 2004; Guthrie & Evans, 2004; Malamud
71 et al., 2004; Dong & Guzzetti, 2005; Van Den Eeckhaut et al., 2007; White et al., 2008;
72 Marques, 2008; Brunetti et al., 2009; Rossi et al., 2010; Barlow et al., 2012). However,
73 establishing a numerical link between the scaling parameters and environmental

74 conditions represents a significant challenge in applying power law statistics to
75 predictive models of landscape evolution (Barlow et al., 2012).

76 The frequency of failure events along coastal cliffs is generally much higher
77 than that of terrestrial cliffs making them ideal for magnitude-frequency analysis. This
78 research assesses the erosion of chalk cliffs at Telscombe, UK using an inventory of
79 mass movements derived from 12 months of high-precision monitoring. The cliffs are
80 predominantly composed of Newhaven Chalk, which is mechanically weak. High-
81 magnitude instability within the chalk is controlled by sliding across two key joint sets
82 such that a wedge-type failure mechanism is common (Mortimore et al., 2004a; Barlow
83 et al., 2017). Wave attack at the base of cliff influences stability both through direct
84 erosion and through microseismic accelerations. Brain et al. (2014) demonstrated that
85 microseismic shaking induced by storm waves can exceed the strength properties of
86 weaker rocks in areas of high stress concentration. More recently, Earlie et al. (2015)
87 observed that high-magnitude waves can cause ground accelerations an order of
88 magnitude greater than have previously been observed and rates of erosion two
89 orders of magnitude larger than the time-averaged mean for cliffs formed of
90 mudstones, siltstones, and sandstones. The shear strength of joints is significantly
91 influenced by the presence of co-planar rock bridges (Bonilla-Sierra et al., 2015).
92 These results suggest that crack nucleation and progressive failure of rock bridges is
93 likely accelerated due to microseismic shaking associated with high-magnitude storm
94 waves. We hypothesise that the weakness of the chalk combined with the structural
95 control and microseismic accelerations associated with high-magnitude storm waves
96 act to minimise the time of progressive failure. Observations over long sections of cliff
97 that include multiple potential failure blocks should therefore provide a quantitative link
98 between mass movements and wave action at the base of cliff. **Indeed, this research**

99 aims to provide a constraint for the power law parameters so that future erosion at
100 Telscombe can be determined using probabilistic erosion models. These models are
101 based on the current and future marine projections provided by the UKCP09 medium
102 emission forecast between 2020 and 2090 and account for sea level rise through a
103 time of exposure approach. Our method is applicable to any coastal cliffs composed
104 of weak rock that respond quickly to basal erosion. Results provide a probabilistic
105 recession model that should be of great use to coastal managers.

106 Study site

107 Telscombe cliffs are located along the southeast coast of the UK as shown in Figure
108 1 and represent one of the few sections of sea cliff between Brighton and Newhaven

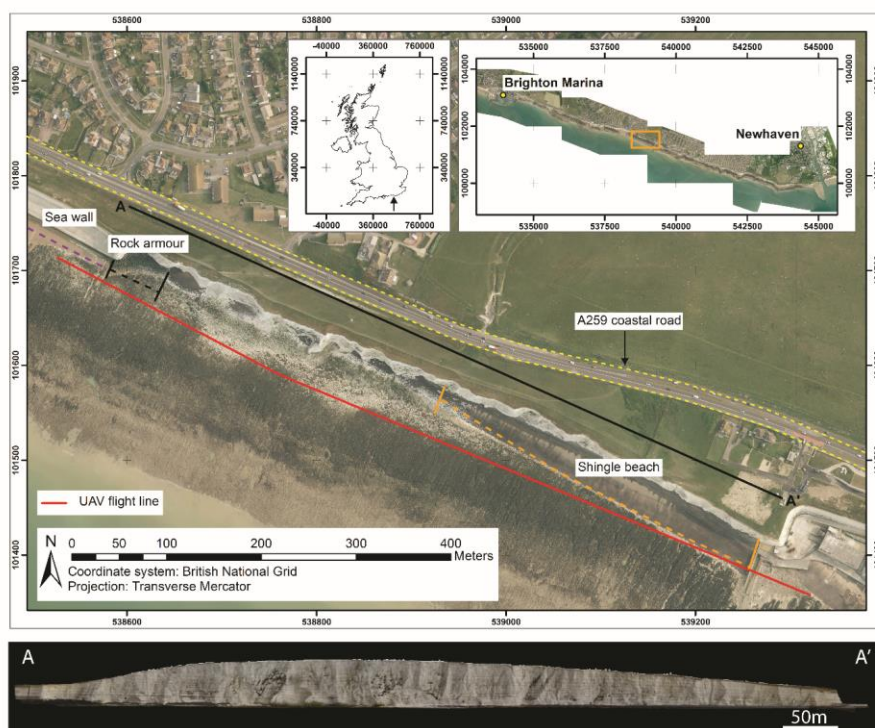


Figure 1: Study area at Telscombe, UK with natural and artificial cliff toe protection identified and cliff model produced through photogrammetry (2013 aerial imagery downloaded from the Channel Coastal Observatory CCO).

109 that remains undefended. The main coastal road (A259) runs within 42.1m of the cliff
110 edge with an average monthly traffic flow of 21,450 vehicles between 2013 and 2016
111 (Brighton & Hove City Council, 2017). Formed of Cretaceous Chalk of the Newhaven
112 and Culver formations, the cliffs are approximately 750 m in length with dry valleys at
113 either end of the site (Mortimore, 1997). The maximum elevation is approximately 49
114 m and the cliffs are orientated towards the dominant wave direction from the south-
115 west (May, 2003). To the west the cliffs have been decoupled from wave action
116 through construction of a sea wall and promenade and have been artificially regraded.
117 Rock armour protects the toe of the cliff for the western 50 m of the site to prevent
118 outflanking of the sea wall (Figure 1). At the eastern extremity a concrete groyne
119 which protects the sewage outfall pipe acts as a barrier to the transport of sediment.
120 As a result, a substantial shingle beach protects the toe of the cliff which tapers over
121 approximately 300 m (Figure 1). The site is macro-tidal with an average spring tidal
122 range of 6.1 m (CCO, 2015) which submerges the shore platform at high tide and
123 enables wave interaction with the base of cliff on a daily basis along the exposed
124 section. Where the cliff is protected by the shingle beach wave interaction occurs
125 during storm events for approximately the western 50-100m when the shingle has
126 been removed from the foreshore. For the remaining section of cliff line the shingle
127 beach remains and therefore wave interaction for this section is minimal. Annual
128 rainfall averages at 720 mm and significant wave heights for this section of coast
129 average at 0.64 m and 1.04 m for the summer and winter months respectively (CCO,
130 2015). Kinematic analysis of the chalk at the site indicates two steeply inclined joint
131 sets as illustrated in Figure 2 ([Barlow et al., 2017](#)). The orientation of these relative to
132 the strike of the cliff face indicates wedge failure as the most likely mode of slope
133 instability. These structural controls produce the characteristic pyramidal morphology

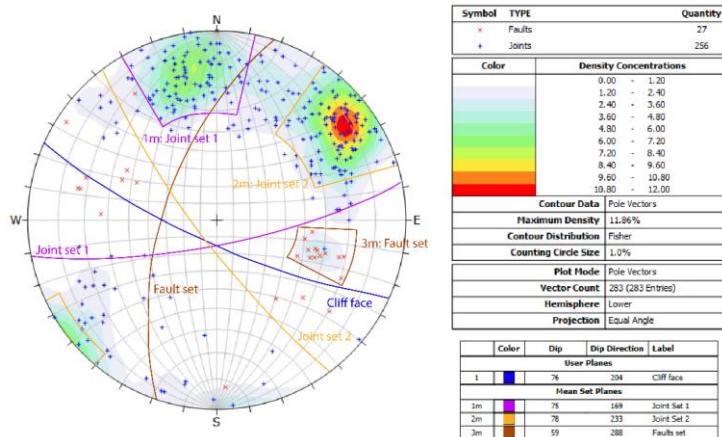


Figure 2: Wolf equal angle stereographic projection of discontinuity data from Telscombe cliffs showing pole plot densities of mapped joints and faults.

of Newhaven formation cliffs (Mortimore et al., 2004a). When saturated, Newhaven Chalk has a cohesion of 600 KN/m² and an angle of internal friction of 24° which increase to 2400 KN/m² and 43.5° respectively when dry (Mortimore et al., 2004b). This weakness suggests that the chalk will respond quickly to changes in the stress environment brought on by wave erosion.

Methods

This study utilises UAV photogrammetry to populate monthly mass movement inventories based on change detection between sequential 3D cliff models. The resultant magnitude-frequency data were used to develop negative power law models whose paramaters s and β were correlated with H_s . The probability of the maximum failure volume in any given month was used to complete the methodology. Each method and parameter is explained in the following sections.

UAV photogrammetry for obtaining a rockfall inventory

High precision monitoring of the site was conducted using a DJI S1000 octocopter, fitted with a Nikon D810 FX DSLR 36 mega-pixel camera with an AF Nikkor 24mm f/2.8D lens. The camera settings were optimised for the aircraft speed; aperture f/8 and shutter speed of 0.002 (1/5000) sec and natural lighting conditions ISO ranged from 800-1600, over the twelve surveys. An automated flight path was used for each survey which maintained a distance of 50 m between cliff face and the camera with a flying altitude of 21 m (approx. mid-cliff height). During the flights, the camera was maintained orthogonal to the cliff face via live streaming onboard video. A traditional strip plan was selected as the best method of capture for long stretches of cliff line captured from a relatively close distance with short focal length (Birch, 2006). Flight speed was set at 3 ms⁻¹ with an image capture interval of 5 s such that an image was captured every 15 m. Initial surveys utilised a ground control network of 23 points located using differential global positioning system and total station surveys. Following the initial surveys, a network of flints was selected and their coordinates extracted from the models to form a relative control network which reduced the field surveying time. Survey accuracy produced a 3D standard error of 0.05 m. Bundle adjustments, generation of epipolar images and point cloud generation was undertaken in the ADAM 3DM software environment. Point clouds which had an average density of 351 points/m² were then rasterised to a cell size of 0.1m before a 2.5D change detection (Rosser et al., 2005) was undertaken between successive datasets. The 2.5D change detection was undertaken using the average plane of the cliff which was recorded at 204°. Sub-sections of the study area were investigated and showed ±4°, the impact on the results would be considered to be negligible, therefore only one plane was used in all subsequent 2.5D change detection analysis. Error assessment found the greatest component residual error from any monthly dataset was 0.10m, as a result volumetric

estimations were calculated with a minimum reliable detectable rockfall size of $1 \times 10^{-3} \text{ m}^3$. A total of 10,085 failures were recorded over the 12 months.

Negative power law parameter estimation

Magnitude-frequency histograms were plotted on logarithmic axes (Figure 3A) using logarithmic binning methods (Guzetti et al., 2002; White et al., 2008; Barlow et al., 2012). Frequency densities were calculated for each bin by using the formula (Malamud et al., 2004):

$$f(V_R) = \frac{\delta N_R}{\delta V_R} \quad (2)$$

where $f(V_R)$ is the frequency density of events with magnitude V_R , δN_R is the number of rockfalls within the specified volume range of δV_R , and δV_R corresponds to the width of the bin. The power law parameters were found using least squares regression (LSR)

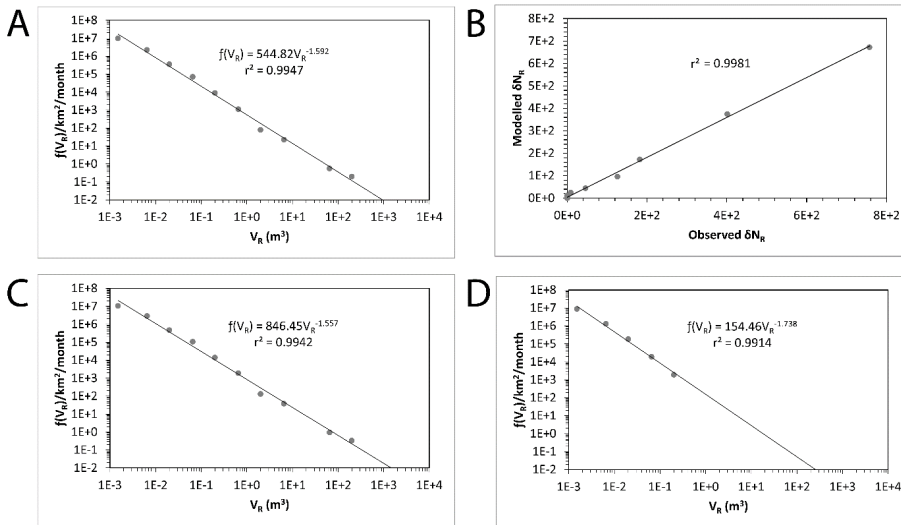


Figure 3: Power law estimation parameters for August to September 2016 (A) frequency density and magnitude of failures for the entire study area, (B) the predicted vs. observed frequency of failures for all binned data, (C) frequency density and magnitude of failures for the undefended section and (D) frequency density and magnitude of failures for the natural defended section (shingle beach) [black lines depict best fit models].

on the logarithmically transformed data (Hovius et al., 1997; Korup, 2005; Barlow et al., 2012). Modelled frequencies were compared to those observed following Barlow et al. (2012) using the integral of Equation 1:

$$\delta N_R = \int_{\min}^{\max} s V_R^{-\beta} dV_R \quad (3)$$

$$\delta N_R = \frac{s V_{R \max}^{1-\beta}}{1-\beta} - \frac{s V_{R \min}^{1-\beta}}{1-\beta} \quad (4)$$

By using Equation 4 and setting the $V_{R \max}$ and $V_{R \min}$ to the bin widths it is possible to assess the accuracy of the estimated power law parameters against the actual observations (Barlow et al., 2012), an example taken from the month August to September 2016 is provided in Figure 3B. The r^2 value of 0.9981 shows definitive agreement between the observations and the power law model providing confidence that modelled frequencies are reliable. As the period between surveys varied, the frequency densities were normalised by time and area ($\text{km}^{-2} \text{month}^{-1}$), with a month represented by 30.4375 days (365.25 days per year / 12 months). The study area was subdivided into 'undefended' and 'naturally defended' to account for the impact of the shingle beach on the power law estimations (E.g. Aug- Sep undefended - Figure 3C & naturally defended – Figure 3D) and subsequent analyses. The r^2 values for the power law estimation parameters (E.g. Figure 3A) varied from 0.9697 to 0.9947 for the entire study area, 0.9693 to 0.9942 for the undefended section and 0.9042 to 0.9979 for the naturally defended section (Table 1).

Month	Entire study area			Undefended (no beach)			Naturally defended (beach)		
	β	s	r^2	β	s	r^2	β	s	r^2
Aug – Sep 2016	1.592	544.82	0.9947	1.557	846.45	0.9942	1.738	154.46	0.9914
Sep – Oct 2016	1.636	175.20	0.9916	1.605	283.91	0.9180	1.546	110.79	0.9879
Oct – Nov 2016	1.573	369.35	0.9984	1.565	613.84	0.9885	1.455	93.79	0.9922
Nov – Dec 2016	1.471	487.50	0.9870	1.453	780.68	0.9868	1.481	167.47	0.9816
Dec – Jan 2017	1.606	265.07	0.9930	1.597	444.46	0.9929	1.452	84.68	0.9979
Jan – Feb 2017	1.645	460.92	0.9875	1.606	782.57	0.9861	1.745	105.46	0.9834
Feb – Mar 2017	1.421	904.14	0.9881	1.409	1459.50	0.9883	1.150	885.55	0.9668
Mar – Apr 2017	1.697	188.83	0.9701	1.682	318.75	0.9693	1.565	53.37	0.9894
Apr – May 2017	1.680	209.69	0.9697	1.669	330.92	0.9724	1.325	255.91	0.9731
May – Jun 2017	1.955	67.76	0.9811	1.829	139.33	0.9805	2.089	28.25	0.9678
Jun – Jul 2017	1.914	33.79	0.9806	1.845	59.36	0.9866	1.300	716.88	0.9042

Table 1: Power law estimation parameters for Telscombe cliffs (Bold identifies those used for the probabilistic modelling).

Modeling cliff erosion with H_s and sea level rise (SLR)

Erosional flux can be calculated for a given magnitude of event by multiplying the frequency density of the event by the magnitude, the result of applying this to the power law equation is (Barlow et al., 2012):

$$V_{RC} = sV_R^{-\beta} V_R \quad (5)$$

$$V_{RC} = sV_R^{-\beta+1} \quad (6)$$

where V_{RC} is the volume in $m^3 km^{-2} month^{-1}$ for an event of magnitude V_R . Thus, the total volumetric erosional flux (V_T) of rock between a minimum and maximum magnitude can be calculated by (Barlow et al., 2012):

$$V_T = \int_{min}^{max} sV_R^{-\beta+1} dV_R \quad (7)$$

$$V_T = \frac{sV_{Rmax}^{-2-\beta}}{2-\beta} - \frac{sV_{Rmin}^{-2-\beta}}{2-\beta} \quad (8)$$

The maximum volume (V_{Rmax}) for Equation 8 can be easily extracted from the rockfall inventory, the minimum value can be more difficult to determine. The minimum detectable rockfall volume using our method was $1 \times 10^{-3} m^3$. To avoid data censoring of smaller volumes (e.g. Stark & Hovius, 2001), a value of $1 \times 10^{-6} m^3$ was used as the minimum threshold in Equation 8.

The distribution of the monthly power law scaling parameters was tested for correlation with H_s , temperature, precipitation, and wind speed data. The atmospheric data was downloaded from the Brighton Marina meteorological station (CCO, 2017a) and the wave data from the distal wave buoy at Seaford, approx. 8.5km from Telscombe at a depth of 11 m CD provided by the Channel Coastal Observatory (CCO, 2017b). To determine the waves which interact with the base of cliff the minimum elevation of the cliff toe was extracted from the photogrammetry models. Tidal data downloaded from the British Oceanographic Centre (BODC, 2017) for Newhaven enabled a binary classification of tidal interaction with the base of cliff. This

229 classification was then applied to the dataset so only waves which interacted with the
230 cliff were considered in the analysis.

231 The projections of future wave climate are presented by the UKCP09 (Lowe et
232 al., 2009) and were generated using the medium emission scenario (IPCC scenario
233 A1B, Leake et al., 2009). This scenario was chosen as it represents the model which
234 global climate change is most closely following (DEFRA, 2009). The wave model was
235 run for the UK continental shelf and meteorological parameters of wind and pressure
236 obtained from the Hadley Centre Met Office were used to force the wave and surge
237 models (Brown et al., 2012). The regional model was divided into approximate 12km
238 grid cells, with the nearest to the study site selected for data extraction. This grid cell
239 was considered an appropriate level of detail for the modelling as it represents a
240 comparable distance between the study site and the distal wave buoy (8.5km). The
241 wave climate is therefore assumed to be representative of the area. A detailed
242 explanation of the wave model and data generation is provided by Brown et al. (2012).
243 The 6-h time series dataset was grouped by month and averaged over 20 years to
244 obtain each decade (e.g. 2010-2029 generates the decade 2020-2029) as is the case
245 with all UKCP09 sub-aerial and marine projections. Probability plots were obtained
246 for each month in each decade between January 2020 and December 2089 (e.g.
247 Figure 4) and a relationship between probability and H_s was found for each month.
248 Random number generation between 0 and 1 could then be used to obtain a H_s value
249 which would calculate β and s using relationships discussed ~~below~~later. To account
250 for SLR the percentage of the tidal cycle that interacts with the base of cliff was
251 recorded, under present conditions 28.58% of the tidal cycle is at or above the
252 minimum cliff height. Using the relative sea level (RSL) rise data presented by the
253 UKCP09 (Lowe et al., 2009), under the medium emission scenario, the 50th percentile

254 (most

likely

scenario)

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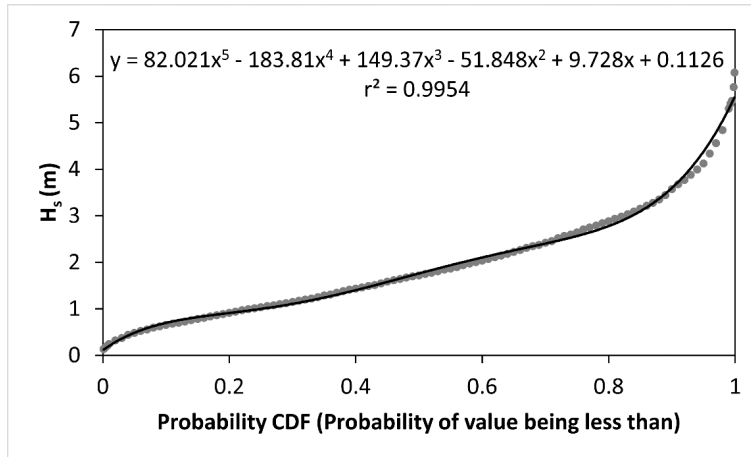


Figure 4: Example of H_s vs. cumulative distribution probability (CDF).

274 extracted for each year. This increase in sea level was added to the observed data
275 and reflected as a percentage of the tidal cycle which interacted with the cliff. From
276 the current observations to 2090 there is a predicted increase from 28.58% to 33.85%.
277 To reflect this in the model the total volume (V_T) calculated from Equation 8 under
278 current conditions whereby 28.58% of the tidal cycle interacts with the cliff was given
279 a scaling factor of 1. The yearly increase in sea level was then given a scaling factor
280 relative to this base level. Therefore, by 2090 when there is a predicted increase of
281 5.27% in the percentage of the tidal cycle interacting with the cliff, the scaling factor
282 applied to V_T for that year was 1.0527.

283 For the $V_{R \max}$ values, a cumulative probability plot was generated (Figure 5)
284 from the inventory so that a random number between 0 and 1 could generate a
285 maximum failure volume to be used in Equation 8. As no single model accurately
286 expressed the relationship between cumulative probability and maximum failure, two
287 logarithmic relationships were selected (intersect when cumulative probability = 0.822).

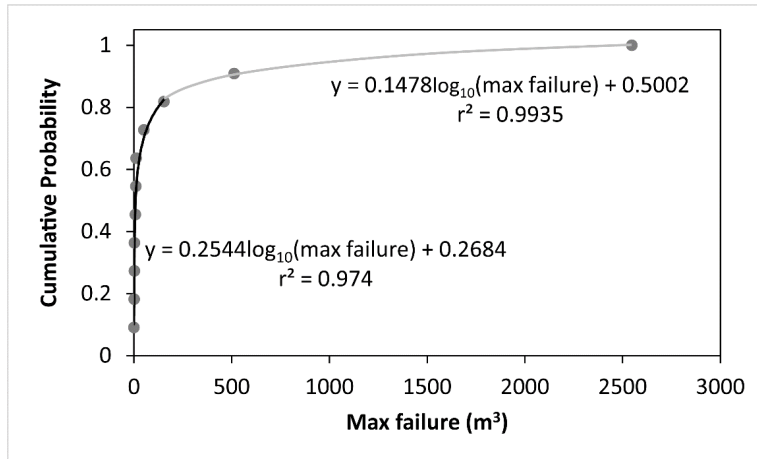


Figure 5: Cumulative probability plot for max failure (black and grey lines indicate best fit models).

288 Using the numerical relationships presented a Monte Carlo simulation model
 289 was run based on Equation 8. To determine the most likely erosion scenario, 10,000
 290 iterations of the model were run from 2020-2090 (70 years). Comparisons between
 291 the model outputs and historical recession rates were undertaken to assess model
 292 reliability. The V_T for 2020-2090 under current conditions, maintained for the entire
 293 temporal period, and future conditions as discussed were then used to calculate cliff
 294 top recession which enabled an assessment of the risk to the A259 coastal road for
 295 this time period.

296 **Model validation**

297 We validated our initial model results by plotting the observed V_T against the modelled
 298 V_T . The r^2 value for the undefended section was very strong at 0.9918 and the model
 299 predicted 97% of the observed V_T for this section. The model performs well with
 300 respect to the months with the largest failures and the model overestimates V_T when
 301 the total volumetric flux for a given month is less than approximately 10 m³ as also
 302 reported by Barlow et al. (2012). The undefended section produced the stronger
 303 model with respect to H_s as the neighbouring section of cliff with the natural beach

304 offered substantial protection from wave action affecting the toe, and therefore the
305 monthly erosion (V_T) of the naturally defended cliff section was significantly reduced
306 (Table 2).

307 Orthorectification of aerial images from 1957, 1973, 1991 and 2013 of the study area
308 was undertaken using photogrammetric methods. Historical recession rates were
309 calculated using the ArcGIS extension - Digital Shoreline Analysis System (DSAS)
310 developed by the United States Geological Society (Thieler et al., 2009). The model
311 predicts average annual recession under current conditions of 0.29m yr^{-1} , the historical
312 recession rate obtained from the orthorectified imagery between 1957 and 2013 for
313 the same area was $0.31\text{m yr}^{-1} \pm 0.16\text{m}$ (2σ) which is similar to other published retreat
314 rates for the site in the scientific literature (Dornbusch et al., 2008). The similarity of
315 the results using different methods of calculating cliff recession provides greater
316 confidence in the model output. The subsequent analysis therefore focusses on the
317 undefended section of cliff line.

318 **Results and discussion**

319 **2.5D surface change detection**

320 Our high-precision monitoring has provided evidence of erosional cycles at the site
321 focused around the conjugate joint sets. Figure 6 illustrates the 2.5D surface change
322 detection results which identify toe erosion between August 2016 (Figure 6A) and
323 February 2017 (Figure 6F) totalling 104.90m^3 , concentrated in the lower section of the
324 cliff. As a result of this erosion, the mean slope angle of the cliff face increased from
325 $\sim 76^\circ$ to $\sim 80.5^\circ$. This steepening resulted in a large wedge failure observed in March
326 2017 (Figure 6G). Over time wave action will remove the debris at the base of cliff
327 and restart the cycle.

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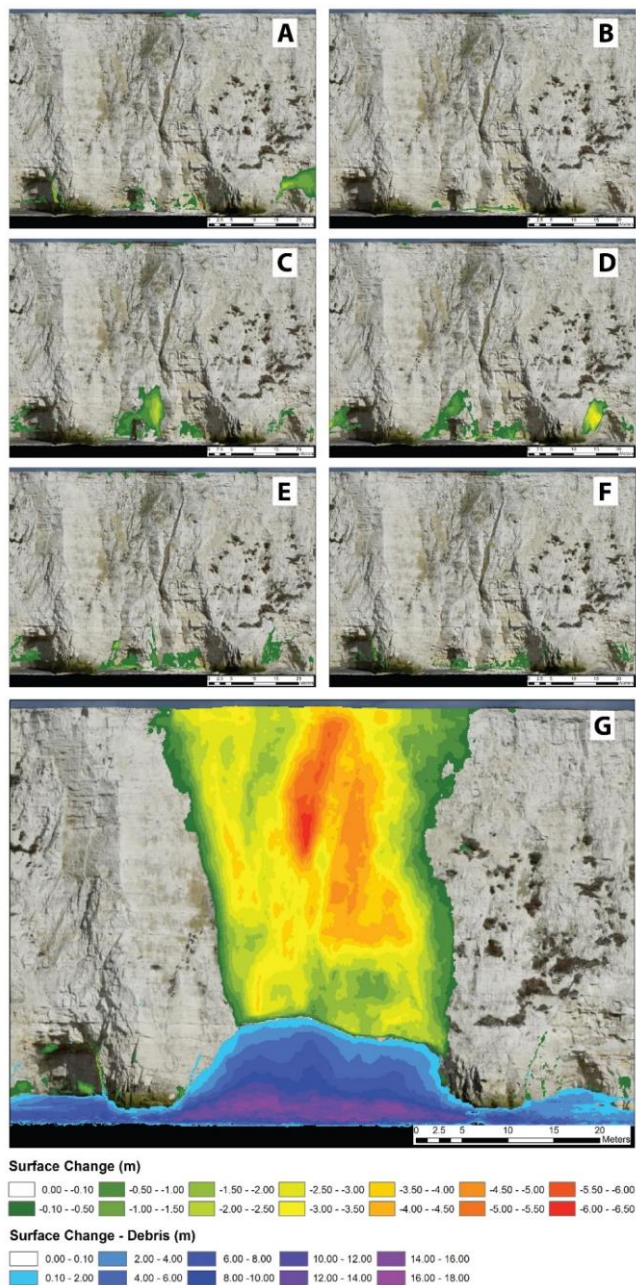


Figure 6: Surface change between August 2016 and March 2017, illustrating a time series of toe erosion and wedge failure (A) Aug-Sep 2016, (B) Sep-Oct 2016, (C) Oct-Nov 2016, (D) Nov-Dec 2016, (E) Dec 2016-Jan 2017, (F) Jan-Feb 2017 and (G) Feb-Mar 2017.

Erosion modelling

We sampled mass movements across 7 orders of magnitude. Given that the largest event in our inventory involved the entire rock face at the highest point in the cliffs, we are confident that we have data covering the entire range of failure magnitudes that are possible at the site. The inventory is comprised of 10,085 mass movements with a total volumetric flux of 3,889.35 m³. The surface area of the cliff face for the entire study area is 23,597.45 m², this was subdivided into the undefended section and the naturally defended beach section with surface areas of 13,890.46 m² and 9,706.99 m² respectively. Although accounting for approximately 59% of the study area, the undefended section had a total volumetric flux of 3,872.62 m³, representing 99.57% of the entire study area total. The monthly inventories are listed in Table 2 and comparisons made between the total volumes (V_T) for the undefended (no beach) and naturally defended (beach) sub-sections. The largest failure was observed between February and March 2017 with an estimated volume of 2,546.81 m³. Other notable failures occurred between November and December 2016 and in the late summer of 2016 (August – September 2016) with estimated volumes of 512.23 m³ and 152.66 m³ respectively. Correlations can be drawn between these larger failures and the distal wave environment of the period between surveys. The average H_s for Aug-Sep, Nov-Dec and Feb-Mar were 1.26 m, 2.12 m and 2.31 m respectively. Furthermore, the percentage of waves which interacted with the cliff that were classified as distal storm waves, $H_s > 3.85$ m (Bovington et al., 2015), were 6% (Aug-Sep), 9% (Nov-Dec) and 14% (Feb-Mar), the highest percentage in the remainder of the dataset was 3% between May and June. These results indicate a minimal lag time between delivery of high energy waves and large magnitude events, either through

353 toe erosion or a significant increase in the stress environment propagating along
 354 conjugate joints (Sunamura, 2015).

Month	V _T Observed (m ³)	
	Undefended (no beach)	Naturally defended (beach)
Aug – Sep 2016	252.59	2.83
Sep – Oct 2016	22.05	1.23
Oct – Nov 2016	119.16	0.81
Nov – Dec 2016	622.01	1.02
Dec – Jan 2017	64.41	0.47
Jan – Feb 2017	31.42	1.83
Feb – Mar 2017	2702.04	1.59
Mar – Apr 2017	20.07	0.67
Apr – May 2017	32.03	3.79
May – Jun 2017	4.59	1.8
Jun – Jul 2017	2.24	0.69

355 Table 2: Monthly observed total volumetric erosion (V_T) for the undefended and naturally defended
 356 cliff line at Telscombe.

357 **Scaling parameters**

358 Significant linear regression correlations, **at the 95.5% (2 σ) confidence level**, were
 359 found between H_s and the β and s values with r^2 values of 0.4971 and 0.5902
 360 respectively for the undefended section of cliff (Figure 7A & 7B). A dataset with a
 361 greater spatial extent would likely increase the strength of these relationships as the
 362 site has a limited number of conjugate joint sets where wedge failures occur, each at
 363 a different stage of the failure cycle. Reoccurring failures are therefore temporally
 364 limited by the erosional cycle. With a greater sample of conjugate joint sets the
 365 likelihood that delivery of high energy waves would lead to failure is increased, thus
 366 strengthening the relationships between H_s and the power law scaling parameters.
 367 The linear regression relationship between s and H_s intersects the x axis at 0.602 m
 368 for the undefended section (Figure 7B). When comparing the erosion activity to the
 369 H_s values from the observation dataset it was considered appropriate to determine
 370 that if the value of H_s in the model was below this threshold the erosion activity would
 371 be set to zero. **A further control site at Brighton marina, where the cliffs are**

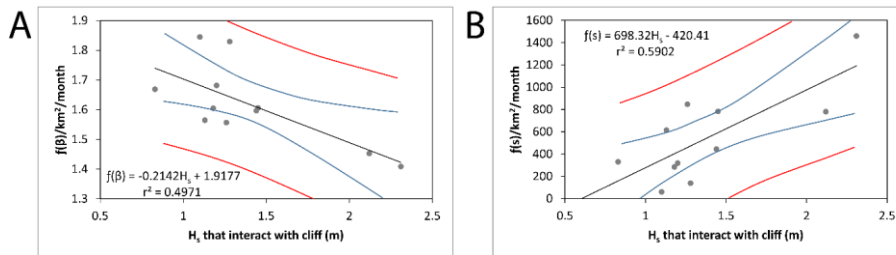


Figure 7: Normalised power law estimation parameters β and s vs. H_s (A) β vs. H_s for the undefended section of cliff and (B) s vs. H_s for the undefended section of cliff (black lines indicate best fit models, blue lines the 95.5% CI and red lines the 95.5% predicted values range).

~~protected-disconnected from wave energy by a sea wall,~~ recorded no detectable erosion during the same period. Although sub aerial and marine erosion below the threshold do occur, the influence of these on the total overall mass flux appears to be likely to be negligible/minimal. No significant correlation was found between the scaling parameters and wave direction (r^2 values of 0.0884 for β and 0.0293 for s), or sub-aerial conditions of precipitation (r^2 values of 0.0002 for β and 0.0006 for s) and air temperature (r^2 values of 0.2628 for β and 0.2178 for s) for the undefended section. Likewise there were no significant correlations found for the naturally defended beach section for any marine or sub-aerial variables considered.

The s value provides an indication on the level of activity within a given dataset whereas β describes the relative contribution of high-magnitude events to the total volume (Barlow et al., 2012). For example, as β increases, smaller magnitude events contribute a greater amount to the total volume than the larger magnitude failures. The normalised scaling parameters for the undefended section of cliff vary between 1.409 to 1.829 for β and from 59.36 to 1459.50 for s . The range in β values is consistent with those presented in other scientific research (Van Den Eeckhaut et al., 2007; Brunetti et al., 2009; Barlow et al., 2012). The smaller β and larger s values are generally found in the winter months (winter ave. β =1.537, ave. s =895.51, summer

ave. $\beta=1.744$, ave. $s=348.38$) when an increase in erosion would be expected due to the frequency of storms. With regards to the maximum failure volume (V_{max}), this is primarily controlled by the slope morphology (Martin et al., 2002). By finding statistically significant relationships between the scaling parameters (β & s) with H_s the power law model can be constrained to provide a useful predictive capability for future coastal management which is presented in the following section.

Erosion scenarios

Results from our Monte Carlo simulations enabled a comparison between predicted erosion under current conditions and that predicted under UKCP09 conditions. Scaling parameters were controlled through the derived relationships with H_s (Figures 6A & 6B) and the maximum failure volume calculated using the cumulative probability relationships found from the inventory (Figure 5).

The results of the Monte Carlo simulation are provided in Figure 8 and Table 3. The mean total recession increases by 1.31 m (approx.+6%) over the 70 year period

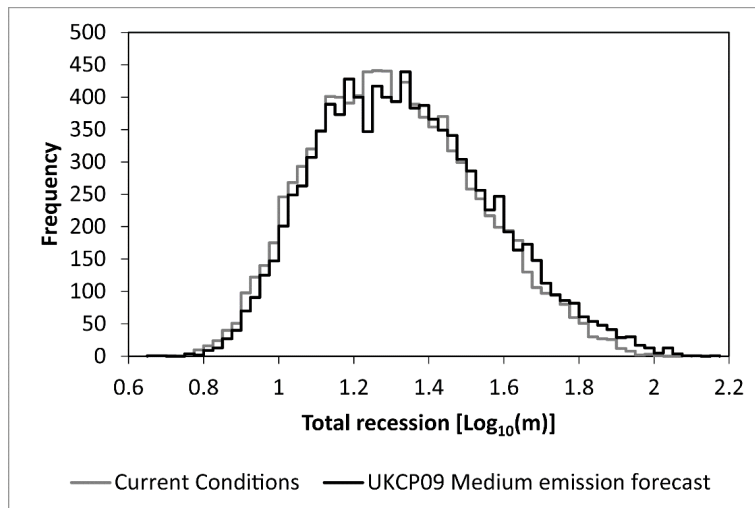


Figure 8: Monte Carlo simulation histogram of Log_{10} transformation of total recession between 2020 and 2090.

404 between the future and current conditions (Table 3), highlighting the impact of future

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425 sea level rise and an increase in significant wave height. The results from the Monte

426 Carlo simulation were normally distributed with a strong positive skew due to the

427 influence of extreme events. In order to apply parametric statistics, the skew was

428 eliminated using a \log_{10} transformation as shown in Figure 8. The maximum recession

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429 distances produced by the model were 121.97 m and 143.56 m under current and
430 future conditions (under the UKCP09 medium emission scenario) respectively. The
431 reason for these extremes values are twofold, firstly due to the high magnitude failure
432 (2546.81 m³) observed in the February to March 2017, this magnitude of failure can
433 be generated in any month of the simulation due to the cumulative probability
434 relationship (Figure 5). This, in combination with projected increases in H_s, enables
435 the model to predict a monthly V_T several orders of magnitude greater than observed
436 in the inventory during the 70 year simulation. However, it should be noted that for
437 these extreme cases the frequency of occurrence was reported as 1/10,000. As the
438 model was run over a 70 year period, the chance of this extreme event occurring is
439 one in 700,000 years. This corresponds to over three standard deviations from the
440 mean under both current and future conditions, reaffirming the unlikelihood of
441 occurrence. The influence of the extreme events impacts the standard deviation
442 resulting in annual recession uncertainty that is approximately three times greater than
443 the calculated historical recession rates (2σ). However, assessing the range of
444 recession values closer (<2σ) to the mean (Figure 8) reveals a significant chance that
445 recession totals between 30 - 60 m over the 70 year period.

	Current conditions		UKCP09 medium emission forecast	
	<u>Recession Log₁₀(m)</u>	Recession (m)	<u>Recession Log₁₀(m)</u>	Recession (m)
Average	1.311	20.45	1.338	21.76
Max	2.086	121.97	2.157	143.56
Min	0.714	5.18	0.659	4.56
95.5%CI	1.750	56.26	1.797	62.61

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446 Table 3: Results of Monte Carlo simulation.

447 The distance between the cliff edge and the A259 ranges between 42.07 m to
448 52.75 m with an average of 46.62 m. Under the most likely erosion scenario (mean in
449 Table 3) approximately 56.6% and 59.7% of the area between the cliff line and
450 minimum distance to the road could be lost by 2090 for current and future conditions
451 respectively (Figure 9). Under both modelled scenarios the 95.5% certainty limit
452 breaches the A259 coastal road (Figure 9) and therefore highlights the potential risk
453 to infrastructure if recession exceeds the modelled average. Assessing the minimum
454 distance from the cliff top to the road approximately 8% of model runs would lead to a
455 breach under current conditions, this increases to about 11% accounting for future
456 conditions as illustrated in Figure 10. The results from this analysis indicate that
457 coastal management policies may need to be reviewed with regards to the current
458 position and future risk to the coastal road. As evidenced with neighbouring sections
459 of cliff line when defences have been installed the annual recession rate has
460 decreased by an order of magnitude. Beach growth is a negative feedback to cliff
461 erosion, where recession results in delivery of flint from the cliff to the beach.

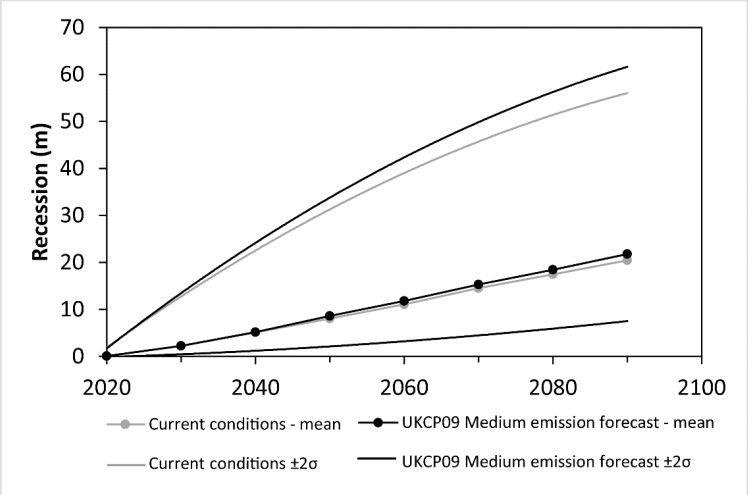


Figure 9: Decadal total cliff recession from Monte Carlo Simulation after 10,000 model runs (Mean recession depicted by data points).

Alternatively accelerated recession within the undefended section could create an embayment and lead to the formation of a pocket beach. The effectiveness of a beach at the cliff toe is evident from the naturally defended section with a significant reduction in erosion (Table 2). The model predictions assume the current state of the site, undefended versus naturally defended (beach), persists into the future.

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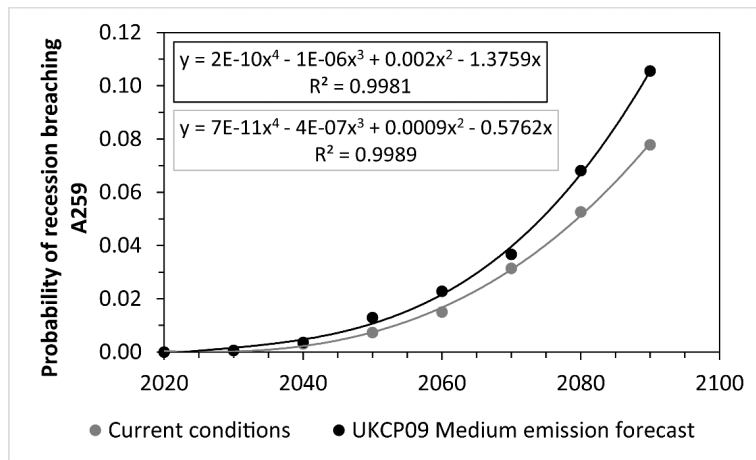


Figure 10: Probability of cliff erosion breaching the A259 at Telscombe cliffs assuming current H_s conditions (grey line) and the UKCP09 forecast for H_s (black line)

Conclusions

This study demonstrates the first statistical link between power law scaling parameters and significant wave height (H_s). ~~This relationship suggests that there is minimal lag between delivery of high energy waves to the base of cliff and failure events. Such relationships offer the potential to quantify the evolution of landscapes under varying environmental boundary conditions.~~ Our results demonstrate the importance of H_s in driving cliff recession at the site ~~and offer the potential to quantify the evolution of landscapes under varying marine conditions. This relationship suggests that there is minimal lag between delivery of high energy waves to the base of cliff and failure events.~~ The simulations involving increased H_s predicted by the UKCP09 model result

477 in a slight acceleration in cliff recession at the site with the most likely outcome being
478 21.7 m by 2090. However, the possibility of extreme events within the model is such
479 that the probability of breaching the A259 increases from ~0% in 2020 to ~11% by
480 2090. The model presented explains ~50% and ~59% of the variation of β and s
481 respectively. Much of the remaining variability is probably controlled by endogenic
482 processes and represents an interesting topic for future research. Although our
483 approach is data hungry, once the inventory data is in place it does not require
484 specialist software to run the simulations such that predictions can be readily updated
485 should improved climate forecasts become available. These methods are
486 transferrable to other sites where the lag between triggering events and mass
487 movement are considered to be minimal. Therefore, our method provides coastal
488 managers with a probabilistic tool to evaluate potential risk to infrastructure through
489 time to facilitate effective planning and mitigation.

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